

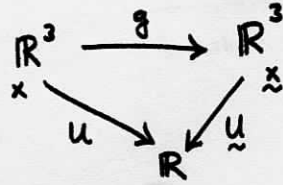
I really should have written
 $G: \mathbb{R}^3 \times \mathbb{R} \rightarrow \mathbb{R}^3 \times \mathbb{R}$
 $(x, t) \mapsto \begin{bmatrix} g(x, t) \\ t \end{bmatrix}$

What is the most general form of co-ord transform
 $g: \mathbb{R}^3 \rightarrow \mathbb{R}^3$ that preserves the form of Newton's
 2nd Law $\vec{F} = m\vec{a}$ [$m\ddot{\alpha}(t) = F(\alpha(t))$] $\frac{d}{dt}(m(t)\dot{\alpha}(t))$

we aren't using the most general form of Newton's Law,
 $= m(t)\ddot{\alpha}(t) + \dot{m}\dot{\alpha}$

(i) Let us consider first the case of a conservative force $F = -\nabla U$.

we know $U = \underline{U} \circ g$
 we want $\nabla_{\underline{U}} = (?) (\nabla U)$



$U = \underline{U} \circ g$

we are ASSUMING m is const.

(a) $\underline{U}(x, t) = U(x, t) = U(g^{-1}(x, t), t)$

Then $\underline{F} = -\nabla \underline{U} = -T(D_x \underline{U}(x, t))$ where T is the transpose operator.
 $= -T(D_x U_x \cdot Dg_x^{-1})$
 $= -(Dg_x^{-1})^T (D_x U_x)^T$
 $= (Dg_x^{-1})^T F$

Later we will show Dg_x is an o.b. matrix.

Then $F(\alpha) = m\ddot{\alpha}$

$\underbrace{D(g^{-1})^T}_\circledast F = m D(g^{-1})^T \ddot{\alpha}$ mult both sides by $(Dg_x^{-1})^T$

$m\ddot{\beta} = \underline{F}(\beta)$

$\Rightarrow \underbrace{m\ddot{\beta}}_\circledast = \underbrace{m D(g_x^{-1})^T}_\circledast \ddot{\alpha}$ (*)

(b) Now we compute $\ddot{\beta}$ from $\beta(t) = g(\alpha(t), t)$ and we see necessary conds on g such that (*) can hold.

This is really the first component of
 $G: \mathbb{R}^3 \times \mathbb{R} \rightarrow \mathbb{R}^3 \times \mathbb{R}$
 $(x, t) \mapsto \begin{bmatrix} g(x, t) \\ t \end{bmatrix}$ $x = \alpha(t)$
 but that does not run this argument

$g: \mathbb{R}^3 \times \mathbb{R} \rightarrow \mathbb{R}^3$ diffeo
 $(x, t) \mapsto$
 $\beta(t) = g(\alpha(t), t)$ transform of trajectory curve of particle.

$\dot{\beta}(t) = D_1 g(\alpha(t), t) (\dot{\alpha}) + D_2 g(\alpha(t), t)$

$\ddot{\beta}(t) = D_1 D_1 g(\alpha(t), t) (\dot{\alpha}, \dot{\alpha}) + \underbrace{D_2 D_1 g(\alpha(t), t)}_{\circledast} (\dot{\alpha}) + D_1 g(\alpha(t), t) (\ddot{\alpha}) + \underbrace{D_1 D_2 g(\alpha(t), t)}_{\circledast} (\dot{\alpha}) + D_2 D_2 g(\alpha(t), t)$

$g: \mathbb{R}^3 \times \mathbb{R} \rightarrow \mathbb{R}^3 \times \mathbb{R}$

cont'd \rightarrow

Thus $\ddot{\beta} = D_1 D_1 g_{(x(t), t)}(\dot{\alpha}, \dot{\alpha}) + 2 D_2 D_1 g_{(x(t), t)}(\dot{\alpha}) + D_1 g_{(x(t), t)}(\ddot{\alpha}) + D_2 D_2 g_{(x(t), t)}$

cut down to 2 components to make writing it out easier.

my notation for this object (a "tensor"?)

$$[\dot{\alpha}_1, \dot{\alpha}_2] \begin{bmatrix} g_{11}^{(0)} & g_{12}^{(0)} \\ g_{21}^{(0)} & g_{22}^{(0)} \end{bmatrix} \begin{bmatrix} g_{11}^{(2)} & g_{12}^{(2)} \\ g_{21}^{(2)} & g_{22}^{(2)} \end{bmatrix} \begin{bmatrix} \dot{\alpha}_1 \\ \dot{\alpha}_2 \end{bmatrix} + 2 \begin{bmatrix} \dot{g}_{11}^{(0)} & \dot{g}_{12}^{(0)} \\ \dot{g}_{21}^{(0)} & \dot{g}_{22}^{(0)} \end{bmatrix} \begin{bmatrix} \dot{\alpha}_1 \\ \dot{\alpha}_2 \end{bmatrix} + \begin{bmatrix} g_{11}^{(0)} & g_{12}^{(0)} \\ g_{21}^{(0)} & g_{22}^{(0)} \end{bmatrix} \begin{bmatrix} \ddot{\alpha}_1 \\ \ddot{\alpha}_2 \end{bmatrix} + \begin{bmatrix} \ddot{g}_{11}^{(0)} \\ \ddot{g}_{22}^{(0)} \end{bmatrix}$$

must vanish vanish This is the only term that can survive vanish

To make $\ddot{\beta} = (D(g_x^{-1}))^T \ddot{\alpha}$ from p. 1, we must have:

(i) $\frac{\partial}{\partial x_j} \frac{\partial}{\partial x_i} g^{(0)} = 0 \quad \forall i, j, k$ [i.e. $D_i^2 g = 0$]

(ii) $\frac{\partial}{\partial x_j} \dot{g}^{(0)} = 0$

(iii) $\ddot{g} = 0$

(iv) AND $D_1 g_x \stackrel{!}{=} (D_1 (g_x^{-1}))^T$

that is to say $(D_1 g_x)^T \stackrel{!}{=} D_1 (g^{-1})_{g(x)}$
 $= (D g_x)^{-1}$

but since $g^{-1} \circ g = Id$
 $D(g^{-1})_{g(x)} D g_x = I$
 $D(g^{-1})_{g(x)} = (D g_x)^{-1}$

And thus $D g_x$ is an O.G. matrix.

▷ so if we Taylor expand g :

$$g^t(x) = g^t(x_0) + D g_0^t(x) + \underbrace{\frac{1}{2!} D^2 g_0^t(x, x)}_0 + \dots$$

using (i)

$$\downarrow \qquad \qquad \downarrow$$

$$h(t) + A(t)x$$

From (ii) $\frac{\partial}{\partial x_j} \dot{g} = 0 \Rightarrow \dot{A}(t) = 0 \Rightarrow A$ does not depend on t

(iii) $\ddot{g} = 0 \Rightarrow \ddot{h} = 0 \Rightarrow h(t) = \vec{a}t + \vec{b}$

$\Rightarrow g^t(x) = \vec{a}t + \vec{b} + Ax$

$D g_x = A$ and from (iv) this is an O.G. matrix.
 This is the most general transform that preserves the form of Newton's 2nd Law

thus $G(x, t) = \begin{bmatrix} g(x, t) \\ t \end{bmatrix} = \begin{bmatrix} Ax + \vec{a}t + \vec{b} \\ t \end{bmatrix}$ "Q" since O.G.

cont'd →

If there is no $U \ni F = -\nabla U$, how do we transform F ?

II Now for a non-conservative F , when we apply a Galilean transform $Q_{a,b,c}(x,t) = Ax + tb + c$, we simply Define $\tilde{F} = D(g^{-1})_{\text{gas}}^T F = AF$

and then the preceding argument goes thru unchanged.

III Show the Galilean transforms form a group: where A is O.G.

$$G_{a,b,c}(x,t) = \begin{bmatrix} Ax + tb + c \\ t \end{bmatrix}$$

(i) closure

$$H_{p,d,e} \circ G_{a,b,c}(x,t) = (P(Qx + tb + c) + td + e, t)$$

$$= (\underbrace{PQ}_{0G} x + t(Pb + H) + Pte, t) \quad \checkmark$$

(ii) associative

$$K_{r,f,g} \circ (H_{p,d,e} \circ G_{a,b,c}) = (RPQx + t(R(Pb + H) + Pbt + tg) +$$

te + d + e)

(iii) \exists id $Q = \begin{bmatrix} I & x \\ & t \end{bmatrix} = G_{I,0,0}$

(iv) \exists inverse $G_{a,b,c}^{-1} = G_{A^{-1}, -A^{-1}b, -A^{-1}c}$

$H_{p,d,e} \rightarrow$
 $P = A^{-1}$
 $d = -A^{-1}b$
 $e = -A^{-1}c$