

Remark: If I can get work out of a system, the Potential Energy is the work that (presumably) had to be put into the system to get it into that configuration.

See my Fowles AM Ch7 write up May 2009

1/20/2001

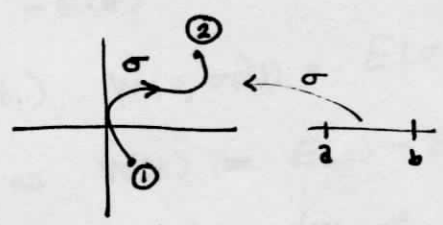
# The Generalized Work - (kinetic) Energy Thm And Conservation of Energy For A "Point Mass"

For systems of particles, see Benson UP p.202

Consider a path  $\sigma$  in the configuration space of a system with a particle of mass  $m$  and a conservative force  $v.f$   $F$  [ $F = -\nabla U$  for some  $U: \mathbb{R}^3 \rightarrow \mathbb{R}$ ]  
Let us assume there is also an external force  $F_e$  acting on the particle.

We want to establish

$$P_1 + K_1 + W_{ext} = P_2 + K_2$$



Consider the following important calculation:

$$\begin{aligned} \text{Work } W_F &= \int_{\sigma} F \cdot ds \\ &= \int_a^b F(\sigma(t)) \cdot \sigma'(t) dt \\ &= \int_a^b \nabla U \cdot ds = - \int_a^b \nabla U(\sigma(t)) \cdot \sigma'(t) dt = U(a) - U(b) = -\Delta(PE) \\ &= \int_a^b F(\sigma(t)) \cdot \sigma'(t) dt \\ &= \frac{m}{2} \int_a^b \frac{d}{dt} (\|\dot{\phi}(t)\|_2^2) dt = \frac{1}{2} m \|\dot{\phi}(b)\|^2 - \frac{1}{2} m \|\dot{\phi}(a)\|^2 = \Delta(KE) \end{aligned}$$

$F \cdot v = \frac{d}{dt} (\frac{1}{2} m v \cdot v)$   
Fowles AM p.86-87

There is a chain of equalities connecting the results of Route A and Route B, thus we have established

$$\begin{aligned} U(a) + \frac{1}{2} m V(a)^2 &= U(b) + \frac{1}{2} m V(b)^2 \\ P_1 + K_1 &= P_2 + K_2 \end{aligned}$$

Total Mechanical Energy is conserved for work done by the conservative force  $F = -\nabla U$

Route B holds for any force  $F$ , not just ones that came from a potential. In particular, if we let  $F_R$  denote the resultant of all forces acting on the particle, we get the Work - (kinetic) Energy Thm (H&R p.99)

$$W_{F_R} = \int_{\sigma} F_R \cdot ds = \frac{m}{2} (V^2(b) - V^2(a)) = \Delta(KE)$$

Thus if there is no change in KE, there is no work done by the resultant force  
H&R p.95, 92

For our system then, we would have  $F_R = F_e + F_u$  [ $F_u$  is just an abbreviation signifying  $F = -\nabla U$ ]

$$\begin{aligned} \text{Thus } \int_{\sigma} (F_e + F_u) \cdot ds &= K_2 - K_1 \\ \Rightarrow \int_{\sigma} F_e \cdot ds &= - \int_{\sigma} F_u \cdot ds + K_2 - K_1 \end{aligned}$$

" $F_e$ " is all the forces on particle, that do NOT come from the potential  $U$ .

$$W_{ext} = P_2 - P_1 + K_2 - K_1 \Rightarrow P_1 + K_1 + W_{ext} = P_2 + K_2 \quad \square$$

Remarks:

- (1) A classic example of this Thm is Bernoulli's Thm in H&R or Benson.
- (2) More generally we would have  $P_1 + K_1 + W_e = P_2 + K_2 + \text{heat} + \text{other energy}$
- (3) Conservation of Energy is the first law of Thermo:  $\Delta E = Q - W$
- (4) Conservation of Energy is restricting the motion (in the phase space) to a co-dim 1 sub mfd.

Masses & Springs Strang ITAM p.42-93

Cont'd ->

## Work-Energy Thm cont'd

We can also establish that power  $P = \frac{dE}{dt}$

For the work done on the system by the external force along an arb path  $\sigma$

$$W_e(a,b) = (P(b) - P(a)) + (K(b) - K(a))$$

$$= \underbrace{[P(b) + K(b)]}_{E(b)} - \underbrace{[P(a) + K(a)]}_{E(a)}$$

$$= E(b) - E(a)$$

Then for any  $t \in (t_0, t_1)$ ,  $W(a, \sigma(t)) = E(\sigma(t)) - E(\sigma(t_0))$

which we can rewrite as  $W(t) = E(t) - E(t_0)$

Then the rate at which work is being done at pt  $\sigma(t)$  and time  $t$

is  $P = W'(t) = E'(t)$

